### Differential Geometry: Definitions and Pictures

#### Nic Ford

May - June 2019

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### What is it?

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We'll be discussing how these work and give some examples of how to use them to do geometry.

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Two natural coordinate systems on half of a sphere.



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$$\operatorname{Hess} f = \begin{pmatrix} 0 & \cos \theta \\ \cos \theta & -r \sin \theta \end{pmatrix}.$$

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...topology prevents you from using the same coordinates for the whole space.

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Even if you only care about  $\mathbb{R}^n$ , many aspects of calculus and geometry appear more natural with the general perspective in mind.

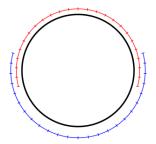
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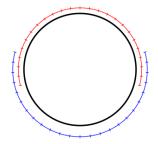
We'll say it consists of a set of points, together with some charts covering all the points. A "chart" is a one-to-one correspondence between some subset of the points and an open ball in some  $\mathbb{R}^n$ .



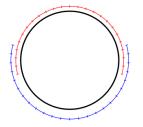
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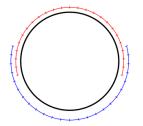


We require the coordinate-change functions from one chart to another to be smooth (i.e. infinitely differentiable).



These two charts can give a circle the structure of a smooth manifold.

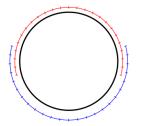
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Writing C for the circle, we can define the charts as:

$$r: \left(-\frac{\pi}{4}, \frac{5\pi}{4}\right) \to C \qquad b: \left(\frac{3\pi}{4}, \frac{9\pi}{4}\right) \to C$$
$$r(t) = (\cos t, \sin t) \qquad b(u) = (\cos u, \sin u)$$

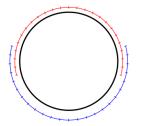


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The red-to-blue coordinate change function  $b^{-1} \circ r$  is defined on  $\left(-\frac{\pi}{4}, \frac{\pi}{4}\right) \cup \left(\frac{3\pi}{4}, \frac{5\pi}{4}\right)$ .

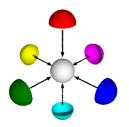


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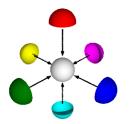
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Six charts on a sphere.



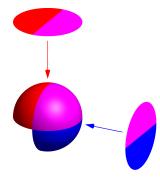


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Writing S for the sphere and B for the open unit ball in  $\mathbb{R}^2$ , two of these charts might be:

 $r: B \to S$   $b: B \to S$ 

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The red-to-blue coordinate change function is

$$(u,v)\mapsto (v,\sqrt{1-u^2-v^2}).$$





#### Four charts on a torus.

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This definition allows you to say which functions on a manifold are smooth, and which maps between manifolds are smooth. (More on this next time!)

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This definition allows you to say which functions on a manifold are smooth, and which maps between manifolds are smooth. (More on this next time!)

The charts give many local coordinate systems that overlap with each other, but no particular one is special.

### Building smooth manifolds

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on  $\mathbb{R}^n$ .

Two ways to build a smooth manifold:

- Define the charts explicitly and give rules for how to glue them together.
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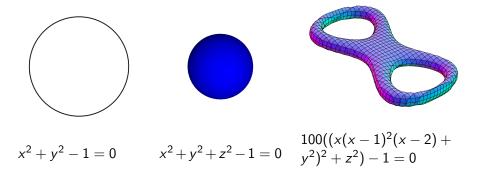
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The implicit function theorem says that, for any system of equations whose Jacobian matrix has full rank at p, we can find coordinates around p where the equations look like this.

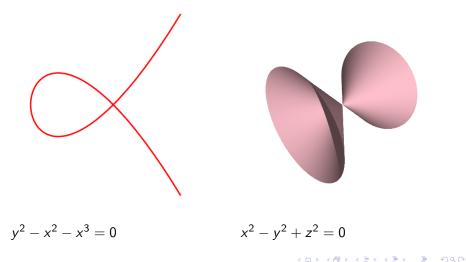
#### Some examples

There are all smooth manifolds defined by one equation each:



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These equations fail to define smooth manifolds because they violate the condition in the implicit function theorem at one point:



The set of all 3D rotations and reflections forms a three-dimensional manifold called O(3).

Each rotation or reflection in O(3) corresponds to a  $3 \times 3$  matrix A for which  $AA^T = 1$ .

These conditions can be written as equations in the entries of the matrix, so O(3) can be cut out by six equations in  $\mathbb{R}^9$ .

The method of defining manifolds as solutions to equations makes it even harder to think in terms of any particular coordinate systems!

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#### Questions?

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# Smooth functions and smooth maps

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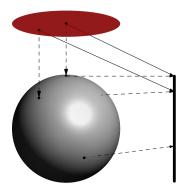
We'll start by asking what it means for one of these to be smooth, and then we'll discuss how to differentiate them.

## Smooth real-valued functions on manifolds

We say a real-valued function  $\phi : M \to \mathbb{R}$  is smooth if, when you restrict it to each chart on M, it's a smooth function on the corresponding open ball in  $\mathbb{R}^n$ .

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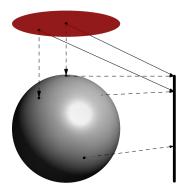


Example: the function on the sphere that takes each point to its *z* coordinate (as point in  $\mathbb{R}^3$  with the usual embedding of the sphere).

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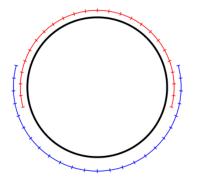
Because we demanded that our coordinate-change functions have to be smooth,  $\phi$  is smooth near some point p if and only if it's smooth on *any* chart containing p.

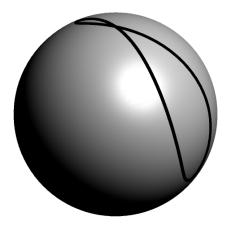
For a map  $f: M \rightarrow N$  from one manifold to another, we again define smoothness in terms of charts:

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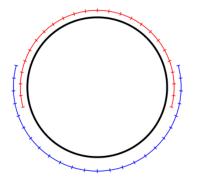
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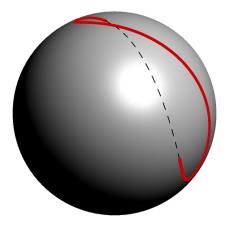
Our map is smooth if, for each  $p \in M$ , there's an M chart around p and an N chart around f(p) on which the corresponding map between open balls in  $\mathbb{R}^n$  is smooth (wherever it's defined).





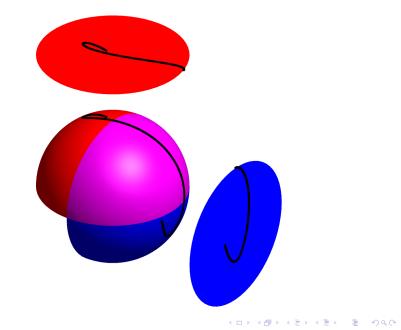
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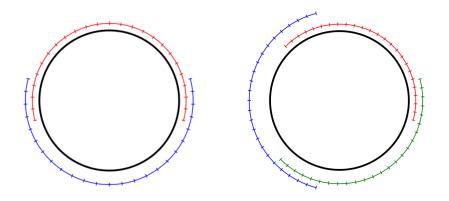
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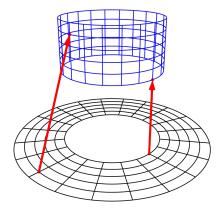
Again, if we get this on *any* collection of charts that cover the image of M, it follows for *every* chart we could check.

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When are two manifolds "the same"?



# Diffeomorphisms

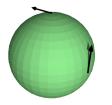


Two manifolds are often thought of as "the same" if they are diffeomorphic: there's a one-to-one correspondence between two the points that's a smooth map in both directions.

Derivatives are linear maps that describe how f(x) changes in response to small changes in x.

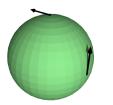
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The input to a derivative should be a direction, but there's no reason for directions at one point on a manifold to have *anything to do with* directions at another point!

So we should attach a *different* vector space to each point of the manifold, called the tangent space.

# Defining tangent spaces

Here are three approaches:

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# Defining tangent spaces

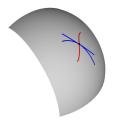
Here are three approaches:



If the manifold is embedded in some  $\mathbb{R}^n$ , we can use the literal tangent space (or line, or plane). This picture is just a little misleading: tangent vectors should be thought of as extremely *close* to the point, not extending far away from it.

# Defining tangent spaces

Here are three approaches:



We can think of tangent vectors as being represented by tiny curves through the point — that is, smooth maps  $\gamma: (-1,1) \rightarrow M$  with  $\gamma(0) = p$ . We declare that two curves represent the same tangent vector as each other if they have the same *velocity*. It's harder to see how to add these, but it can be done.

We can simply *define* the tangent vector in terms of the derivative operator.

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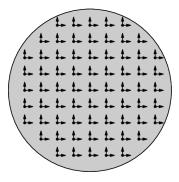
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We'll write  $T_x M$  for the tangent space to M at the point x (i.e., the vector space containing all the tangent vectors).

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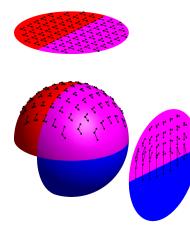
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Picking a coordinate system on a chart gives rise to a collection of tangent vectors at every point in the chart called coordinate tangent vectors. At each point x, the coordinate tangent vectors for a basis for  $T_x M$ .

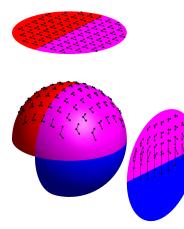
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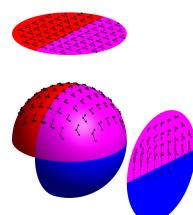


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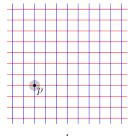


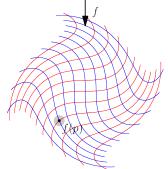
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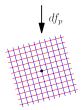
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The relationship between the coordinate vectors works out to:  $\hat{u} = (-\sqrt{1-s^2-t^2}/t)\hat{t}; \ \hat{v} = \hat{s} - (s/t)\hat{t}.$  Any smooth map  $f: M \to N$  gives us a *linear* map  $df_p: T_pM \to T_{f(p)}N$  on all of the tangent spaces called the pushforward or total derivative. (You might see the notation  $f_*$  too.) Think of it as "what f looks like if you zoom in very close to x."









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When you pick charts around p and f(p), the total derivative becomes the ordinary Jacobian matrix.

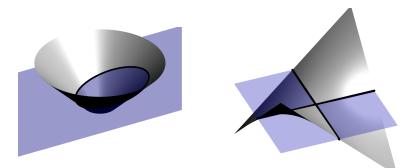
$$\begin{pmatrix} \frac{\partial f_1}{\partial x_1} & \cdots & \frac{\partial f_1}{\partial x_n} \\ \vdots & \ddots & \vdots \\ \frac{\partial f_m}{\partial x_1} & \cdots & \frac{\partial f_m}{\partial x_n} \end{pmatrix}$$

### The tangent space as the "zoomed-in" manifold

In many cases, the tangent space to a manifold contains the relevant information about the behavior of manifold around the corresponding point.

# The tangent space as the "zoomed-in" manifold

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Example: if the intersection of two submanifolds is transverse — meaning their tangent spaces always intersect in the lowest dimension possible — the result is another manifold.

Consider a manifold M embedded in  $\mathbb{R}^n$  via a map  $i : M \to \mathbb{R}^n$ . The tangent plane (in the ordinary sense) at a point p is the image of  $di_p$ .

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From this perspective, the chain rule is just the fact that total derivatives respect composition of functions, that is,  $d(g \circ f)_x = dg_{f(x)} \circ df_x.$ 

#### Questions?

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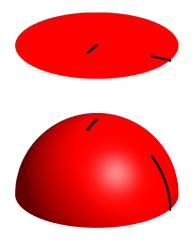




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Just picking coordinates is no good: nothing requires our coordinate charts to preserve lengths! Most of the time they don't and shouldn't be expected to.

We need to put an extra piece of structure on our manifold to do this. The thing we'll use is called a Riemannian metric.



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This is a way of assigning, to each pair of tangent vectors v, w in the same tangent space, a real number  $\langle v, w \rangle$  so that:

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• (positive definite)  $\langle v, v \rangle \ge 0$ , with equality only if v = 0.

An inner product can be thought of as telling you about lengths and angles exactly like the ordinary dot product on  $\mathbb{R}^n$ :

$$\langle \mathbf{v}, \mathbf{w} \rangle = |\mathbf{v}| \cdot |\mathbf{w}| \cdot \cos \theta$$

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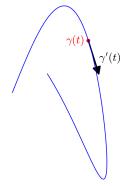
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NB: The inner product is how we *define* lengths and angles! The right-hand sides of these two expressions should be taken as a way of interpreting a choice of inner product as an assignment of lengths and angles to vectors. But the symbols "| |" and " $\theta$ " have no meaning in isolation!

Once we've picked a metric, we can, for example, compute lengths of curves. Given a smooth map  $\gamma : [0,1] \rightarrow M$ , write  $\gamma'(t)$  for the pushforward of the unit tangent vector at t. Then the length of  $\gamma$  is

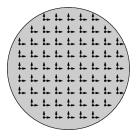
$$\int_0^1 \sqrt{\langle \gamma'(t), \gamma'(t) \rangle} dt.$$



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#### Metrics in coordinates

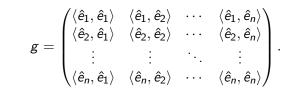
Picking a coordinate chart gives us a basis of coordinate tangent vectors everywhere on the chart. If the coordinates are  $x_1, \ldots, x_n$ , write  $\hat{e}_1, \ldots, \hat{e}_n$  for the coordinate tangent vectors. Once we've done this, we can write our metric as a matrix:



$$g = egin{pmatrix} \langle \hat{e}_1, \hat{e}_1 
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angle & \cdots & \langle \hat{e}_1, \hat{e}_n 
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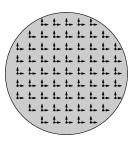
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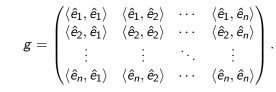
As usual in linear algebra, we can then write  $\langle v, w \rangle = v^T g w$ , provided we expand v and w in the  $\hat{e}_i$  basis.

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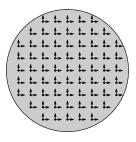
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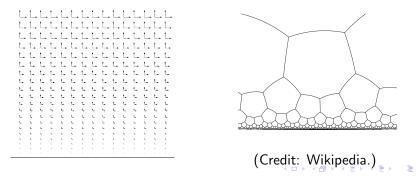
Note that the entries of the matrix will in general be *functions* of  $x_1, \ldots, x_n$ .



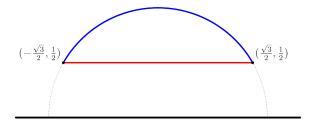
The hyperbolic plane can be thought of as the upper-half plane with the inner product given by

$$\begin{pmatrix} 1/y^2 & 0 \\ 0 & 1/y^2 \end{pmatrix}$$

in the standard basis of coordinate tangent vectors. (What does this mean about the *lengths* of the tangent vectors?)

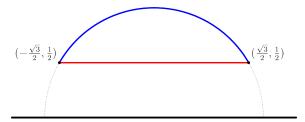


Let's compute the lengths of these two curves:



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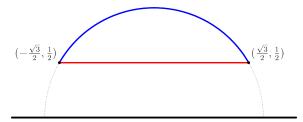
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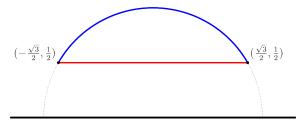


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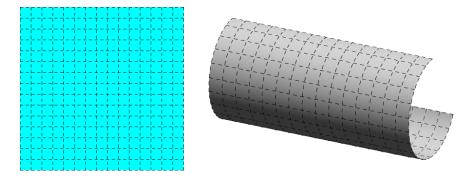
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$$\int_{\frac{\pi}{6}}^{\frac{5\pi}{6}} \frac{1}{\sin t} \, dt = \log\left(\frac{2+\sqrt{3}}{2-\sqrt{3}}\right) \approx 2.634.$$

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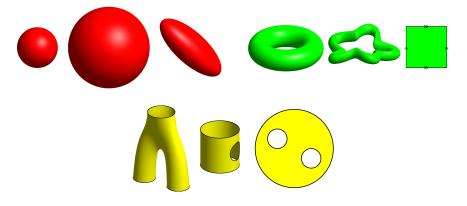
#### Isometries

The presence of a metric introduces a new way for manifolds to be equivalent that's stronger than diffeomorphism: a diffeomorphism  $f: M \to N$  is an isometry if it preserves the metric, that is, we always have  $\langle v, w \rangle_{T_pM} = \langle df_p(v), df_p(w) \rangle_{T_{f(p)}N}$ .



## Isometries

Each group consists of manifolds that are diffeomorphic to each other but not isometric.



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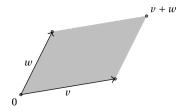
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If you have a coordinate system, you can use it to identify all the tangent spaces with  $\mathbb{R}^n$ , which gives you an obvious choice of metric. But most metrics don't arise in this way for *any* choice of coordinates! The ones that do are called flat. (The usual metric on the sphere, for example, isn't flat.)

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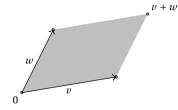


Area<sup>2</sup> =  $\langle v, v \rangle \langle w, w \rangle - \langle v, w \rangle^2$ . Integrate (the square root of) this to get the area of a surface.

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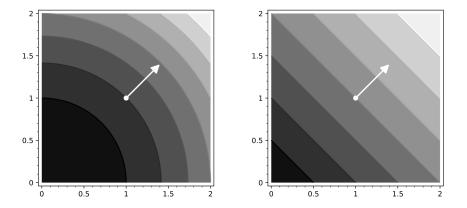
Area<sup>2</sup> =  $\langle v, v \rangle \langle w, w \rangle - \langle v, w \rangle^2$ . Integrate (the square root of) this to get the area of a surface.

In fact, knowing just the lengths is enough to determine the angles, provided that the lengths come from an inner product:

$$\langle v + w, v + w \rangle - \langle v - w, v - w \rangle = 4 \langle v, w \rangle.$$

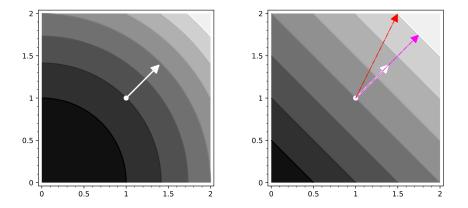
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Some suggestions for things to read to learn more:

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- General Relativity by Robert Wald and Gravitation by Misner, Thorne, and Wheeler. Cover much of the same material motivated by physics (and then also cover the physics).